

# Dinuclear Metalloradicals Featuring Unsupported Metal–Metal Bonds\*\*

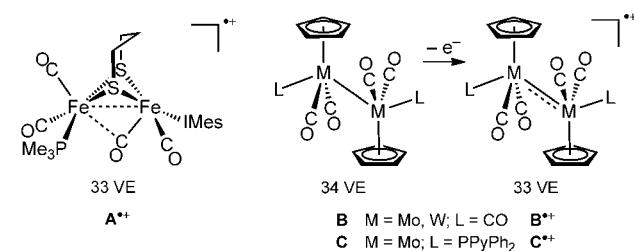
Edwin F. van der Eide, Ping Yang,\* Eric D. Walter, Tianbiao Liu, and R. Morris Bullock\*

Metal–metal interactions in paramagnetic, multinuclear transition metal complexes are critical to the reactivity of metalloproteins,<sup>[1]</sup> and understanding them is important in the development of functional metal-containing polymers.<sup>[2]</sup> Dinuclear mixed-valent complexes—usually obtained by redox reactions<sup>[3]</sup> of diamagnetic precursors—provide insight into such interactions.<sup>[4,5]</sup> Bridging ligands are ubiquitous structural motifs in mixed-valent dimers: dithiolate bridges are found in the active site of [FeFe] hydrogenases and models<sup>[1]</sup> (see **A**<sup>+</sup>, Scheme 1), while bridges formed by unsaturated (hydro)carbon chains and rings,<sup>[5]</sup> carboxylates,<sup>[6]</sup> and amidinates<sup>[7]</sup> are important in other synthetic dimers. Bridging ligands not only mediate the electronic communication<sup>[5]</sup> between the metal centers; they also provide

stabilization against fragmentation. Consequently, much less is known<sup>[8]</sup> about odd-electron dimers that lack bridges and that are only held together by metal–metal bonds.

Group 6 dimers  $[\{\text{CpM}(\text{CO})_3\}_2]$  (**B**, Scheme 1) have unsupported M–M single bonds,<sup>[9]</sup> and their electrochemistry underscores the need for bridging ligands in stabilizing odd-electron dimers. Electrochemical oxidations of **B** are irreversible due to fast decomposition of **B**<sup>+</sup> by cleavage of the M–M bond.<sup>[10]</sup> Derivative  $[\{\text{CpMo}(\text{CO})_2(\text{PPyPh}_2)\}_2]$  (**C**) is reversibly oxidized at low scan rates, but product **C**<sup>+</sup> is still too short-lived ( $t_{1/2} \approx 8$  s)<sup>[11]</sup> to permit its full characterization. The bonding and spectroscopic properties of these interesting unbridged paramagnetic dimers have thus remained unexplored.<sup>[12]</sup> Considering that the M–M formal bond order has been proposed to increase from 1 to 1½ upon oxidation of **B**, **C**, and other group 6 M<sup>I</sup>M<sup>I</sup> dimers,<sup>[12]</sup> we reasoned that stable derivatives of **B**<sup>+</sup> should be accessible. We report herein the first such derivatives that are stable enough to be isolated and fully characterized.

We attribute the high reactivity of radical cation **B**<sup>+</sup> largely to its electron poverty—the higher stability of phosphane-containing **C**<sup>+</sup> confirms this. Hence, usage of the stronger donor phosphane  $\text{PMe}_3$  should further increase the lifetime of the radicals. We synthesized  $[\{\text{CpW}(\text{CO})_2(\text{PMe}_3)\}_2]$  (**1**, see Scheme 2)<sup>[13]</sup> and found that it undergoes an electrochemical oxidation ( $E_{1/2} = -0.43$  V vs  $[\text{Cp}_2\text{Fe}]^+ / [\text{Cp}_2\text{Fe}]$ ,  $\text{CH}_2\text{Cl}_2$  solvent, 50 mM  $n\text{Bu}_4\text{N}^+\text{B}(\text{C}_6\text{F}_5)_4^-$  electrolyte) that remains reversible ( $i_c/i_a > 0.9$ ) at scan rates as low as  $50 \text{ mVs}^{-1}$ . The separation between the anodic and cathodic peak potentials is 65 mV, indicating that one electron is transferred and that  $[\{\text{CpW}(\text{CO})_2(\text{PMe}_3)\}_2]^+$  (**1**<sup>+</sup>) is the oxidation product. Reversibility of the couple **1**<sup>+</sup>/**1** is maintained when the electrolytes  $n\text{Bu}_4\text{N}^+\text{PF}_6^-$  or  $n\text{Bu}_4\text{N}^+\text{BF}_4^-$  (0.1 M) are employed, showing that **1**<sup>+</sup> is (at least on the timescale of the CV experiment) not subject to decomposition by  $\text{PF}_6^-$  or  $\text{BF}_4^-$  anions.<sup>[14]</sup> On a preparative scale, **1**<sup>+</sup> was synthesized as its  $\text{B}(\text{C}_6\text{F}_5)_4^-$  salt<sup>[15]</sup> by reaction of **1** with 1 equiv  $\text{Ph}_3\text{C}^+\text{B}(\text{C}_6\text{F}_5)_4^-$  in  $\text{CH}_2\text{Cl}_2$  (Scheme 2). Crystalliza-



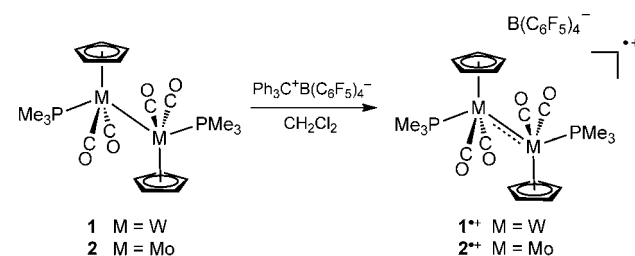
**Scheme 1.** Are bridging ligands required to stabilize paramagnetic dimers? IMes = 1,3-dimesitylimidazol-2-ylidene, Py = 2-pyridyl, VE = valence electron.

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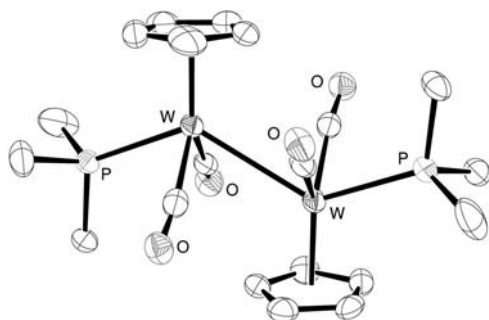
Supporting information for this article is available on the WWW under <http://dx.doi.org/10.1002/ange.201203531>.



**Scheme 2.** Synthesis of the dinuclear metalloradicals **1**<sup>+</sup> and **2**<sup>+</sup>.

tion from  $\text{CH}_2\text{Cl}_2$ /hexanes at room temperature gave  $[\mathbf{1}^+]\text{-B}(\text{C}_6\text{F}_5)_4^{-1/2}\text{CH}_2\text{Cl}_2$  as analytically pure, air-stable, amber crystals in 90 % yield. The dimolybdenum analogue  $\mathbf{2}^+$  ( $E_{1/2}(\mathbf{2}^+/\mathbf{2}) = -0.35\text{ V}$  vs  $[\text{Cp}_2\text{Fe}]^+ / [\text{Cp}_2\text{Fe}]$ ) was similarly accessible, but it was crystallized at  $-35^\circ\text{C}$  to prevent decomposition.

Single crystals of  $\mathbf{1}$  and of  $[\mathbf{1}^+]\text{B}(\text{C}_6\text{F}_5)_4^{-1/2}\text{CH}_2\text{Cl}_2$  were analyzed<sup>[16]</sup> by X-ray diffraction (Figure 1, Table 1). The



**Figure 1.** Structure of  $[\mathbf{1}^+]\text{B}(\text{C}_6\text{F}_5)_4^{-1/2}\text{CH}_2\text{Cl}_2$  (50% ellipsoids; hydrogen atoms, anion, and solvent omitted).

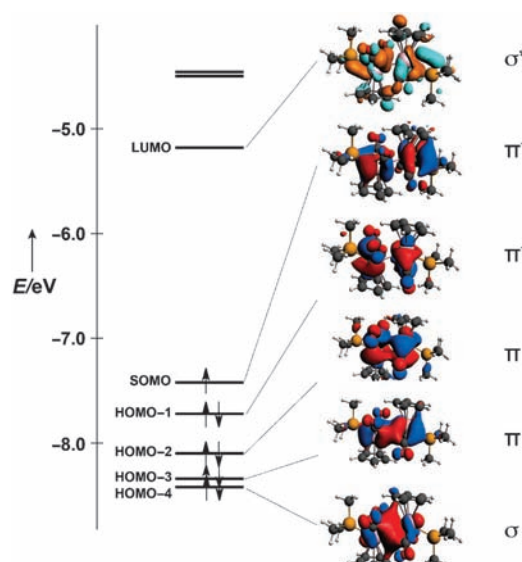
**Table 1:** Selected experimental/calculated bond lengths  $d$  [Å] of  $\mathbf{1}$  and  $\mathbf{1}^+$ .

|                                         | $\mathbf{1}$ [exp] <sup>[a]</sup> | $\mathbf{1}^+$ [exp] <sup>[b]</sup> | $\mathbf{1}$ [DFT] | $\mathbf{1}^+$ [DFT] |
|-----------------------------------------|-----------------------------------|-------------------------------------|--------------------|----------------------|
| W–W                                     | 3.2327(4)                         | 3.0257(3)                           | 3.335              | 3.127                |
| (W–P) <sub>average</sub>                | 2.413(1)                          | 2.468(1)                            | 2.437              | 2.512                |
| (W–C <sub>CO</sub> ) <sub>average</sub> | 1.950(6)                          | 1.968(6)                            | 1.961              | 1.976                |
| (C=O) <sub>average</sub>                | 1.172(7)                          | 1.159(7)                            | 1.186              | 1.177                |

[a]  $C_i$  symmetry (crystallographically imposed). [b] Near- $C_{2h}$  symmetry.

metric differences between  $\mathbf{1}$  and  $\mathbf{1}^+$  are subtle; therefore, only the structure of  $\mathbf{1}^+$  is displayed here (see the Supporting Information for a structure overlay). Both dimers assume the *anti* conformation in the solid state and have W–W bonds that are not supported by bridging or semi-bridging<sup>[17]</sup> carbonyl ligands. The W–W bond length of 3.2327(4) Å in  $\mathbf{1}$  is similar to the length of 3.222(1) Å found in  $[\text{CpW}(\text{CO})_3]_2$ .<sup>[9]</sup> The W–W bond length in  $\mathbf{1}^+$  is 3.0257(3) Å, corresponding to a significant shortening<sup>[18]</sup> (6.4 %) relative to  $\mathbf{1}$  and strongly suggesting that an increase in the W–W formal bond order occurs upon oxidation.

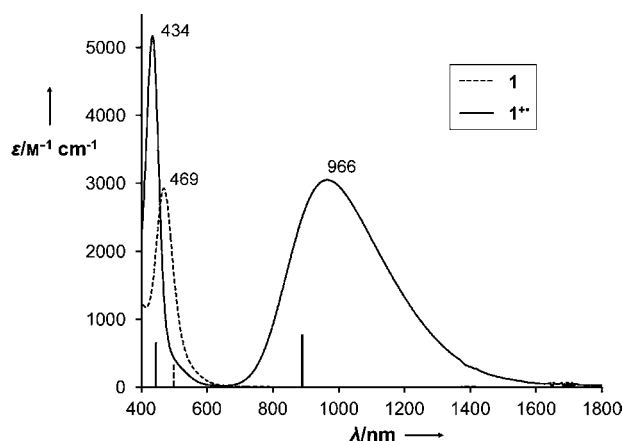
Density functional (DFT; see the *Experimental Section* for details) calculations provided insight into the electronic structures and aided spectroscopic assignments. The optimized structures of  $\mathbf{1}$  and  $\mathbf{1}^+$  are in good agreement with the experimental data (Table 1); the bond-length changes are faithfully reproduced. The MO diagram (Figure 2) clearly shows that the singly occupied MO (SOMO) of  $\mathbf{1}^+$ , which corresponds to the highest occupied MO (HOMO) of  $\mathbf{1}$ , is antibonding and has  $\pi$  symmetry with respect to the W–W bond. Thus, removal of an electron increases the net W–W formal bond order from 1 ( $\sigma$ ) in  $\mathbf{1}$  to  $1\frac{1}{2}$  ( $1\sigma + \frac{1}{2}\pi$ ) in  $\mathbf{1}^+$ . The SOMOs of  $\mathbf{1}^+$  and  $\mathbf{2}^+$  have largely metal character; spin densities of about 95 % (divided equally over both metal centers) are calculated for the M–M cores.



**Figure 2.** Frontier molecular orbital (MO) energy levels of  $\mathbf{1}^+$  and plots of orbitals with metal d character.

The IR spectrum of  $\mathbf{1}^+$  in  $\text{CH}_2\text{Cl}_2$  features one broad carbonyl stretching ( $\nu_{\text{CO}}$ ) band at  $1875\text{ cm}^{-1}$ .<sup>[19]</sup> Although two bands ( $A_u$  and  $B_u$  modes) are expected for a  $C_{2h}$ -symmetric species, the small separation of the bands calculated by DFT [ $(1873 + 1865)\text{ cm}^{-1}$ ] suggests that they are not resolved experimentally. Precursor  $\mathbf{1}$  has bands at  $\bar{\nu}_{\text{CO}} = (1837 + 1809)\text{ cm}^{-1}$  [DFT:  $(1840 + 1817)\text{ cm}^{-1}$ ]. The average  $\bar{\nu}_{\text{CO}}$  therefore increases by about  $50\text{ cm}^{-1}$  upon oxidation,<sup>[20]</sup> which correlates well with an increase in the formal oxidation state of each W atom by half a unit.<sup>[8,21]</sup> The charge (and, therefore, the unpaired electron) in  $\mathbf{1}^+$  is shared equally by both W atoms, even on a very short (ps) timescale, in accordance with the calculations. Any contribution from a localized (i.e.  $\text{W}^{\text{I}}\text{--W}^{\text{II}}$ ) structure can be ruled out.

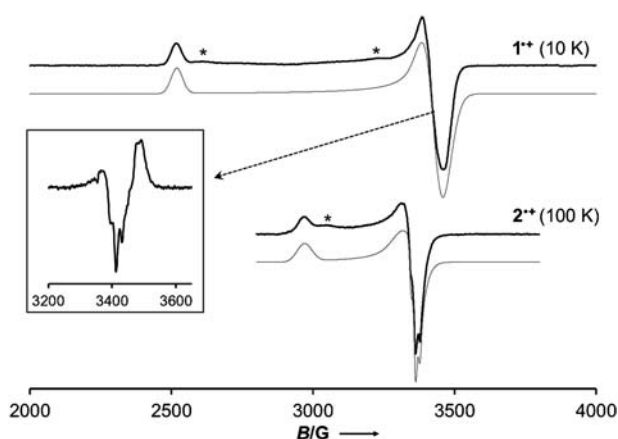
All of the dimers absorb visible light; the W dimers  $\mathbf{1}$  and  $\mathbf{1}^+$  have absorption maxima at  $\lambda_{\text{max}} = 469$  and  $434\text{ nm}$ , respectively (Figure 3). Time-dependent DFT predicts these



**Figure 3.** Vis/NIR spectra of  $\mathbf{1}$  and  $\mathbf{1}^+$  in  $\text{CH}_2\text{Cl}_2$ . The vertical bars mark the transitions predicted by time-dependent DFT calculations, their heights being proportional to the calculated oscillator strengths.

excitations to occur at 499 and 444 nm, respectively, and allows for their assignment as  $\pi^* \rightarrow \sigma^*$  (see Figure 2, HOMO-1  $\rightarrow$  LUMO).<sup>[22]</sup> The open-shell dimers also have characteristic near-infrared (NIR) absorptions at  $\lambda_{\text{max}} = 966$  ( $\mathbf{1}^+$ , Figure 3) and 1110 nm ( $\mathbf{2}^+$ ). Calculations allow us to assign these as  $\pi \rightarrow \pi^*$  transitions involving the SOMO (see Figure 2, HOMO-3  $\rightarrow$  SOMO). NIR transitions are characteristic for mixed-valent compounds and are often described as intervalence charge transfers.<sup>[23]</sup> However, this description implies an electron transfer between metal sites in different oxidation states, which is not applicable to the fully delocalized systems reported herein.

X-band (9.4 GHz) EPR spectra of  $\mathbf{1}^+$  and  $\mathbf{2}^+$  were recorded in frozen toluene solutions (Figure 4). The spectra



**Figure 4.** First-derivative-mode EPR spectra in toluene. Simulations in grey. Inset: derivative of the high-field signal of  $\mathbf{1}^+$ . Resonances marked with an asterisk are assigned to *gauche* isomers.

appear axial as a result of near coincidence of the two lowest principal  $g$  values. However, simulations and calculations (Table 2) reveal the spectra to be rhombic, as expected for

**Table 2:** Experimental/calculated EPR parameters of  $\mathbf{1}^+$  and  $\mathbf{2}^+$ .

|                               | $\mathbf{1}^+$ [exp] | $\mathbf{1}^+$ [DFT] | $\mathbf{2}^+$ [exp] | $\mathbf{2}^+$ [DFT] |
|-------------------------------|----------------------|----------------------|----------------------|----------------------|
| $g_1$                         | 2.662                | 2.606                | 2.258                | 2.229                |
| $g_2$                         | 1.955                | 1.988                | 2.002                | 2.005                |
| $g_3$                         | 1.940                | 1.954                | 1.996                | 1.996                |
| $A(^{31}\text{P})^{\text{a}}$ | 50 <sup>[b]</sup>    | 43.4 <sup>[c]</sup>  | 48 <sup>[b]</sup>    | 42.6 <sup>[c]</sup>  |

[a] In MHz. [b] Observed in one signal. [c] See the Supporting Information.

species of  $C_{2h}$  symmetry. The excellent agreement between calculated and experimental  $g$  values is noteworthy. The high  $g_1$  values and the large  $g$  anisotropies ( $g_1 - g_3$ : 0.72 for  $\mathbf{1}^+$ ; 0.26 for  $\mathbf{2}^+$ ) are good evidence that the spin density is largely metal-based. However, hyperfine coupling to  $^{183}\text{W}$  ( $I = 1/2$ , 14.4 %) or  $^{95}\text{Mo}/^{97}\text{Mo}$  ( $I = 5/2$ , 25 % total) was not discernible due to line broadening. Coupling to  $^{31}\text{P}$  [ $I = 1/2$ , 100 %,  $a_{\text{P}} \approx 17$  G (48 MHz)] is observable in the high-field signal of  $\mathbf{2}^+$ , and becomes apparent in  $\mathbf{1}^+$  when the derivative of the spectrum is shown (Figure 4, inset). The fact that coupling to two equivalent  $^{31}\text{P}$  nuclei is observed in both  $\mathbf{1}^+$  and  $\mathbf{2}^+$  again

emphasizes the symmetric delocalization of the unpaired electron over the M–M cores.

In conclusion, our results show that dinuclear metal-loradicals with no M–M bridging ligands are not inherently unstable. In fact, oxidation of  $[\{\text{CpM}(\text{CO})_2(\text{PMe}_3)_2\}_2]$  leads to dinuclear metalloradicals that have additional  $\pi$  bonding between the metal centers. The increased electron density provided by the  $\text{PMe}_3$  ligands seems to be crucial in stabilizing the radical cations.

## Experimental Section

Synthetic details and characterization data, as well as more complete computational details, are given in the Supporting Information. First-principle calculations based on spin-polarized DFT were performed using the PBE exchange-correlation functional<sup>[24]</sup> implemented in the Amsterdam Density Functional (ADF 2010.02) program.<sup>[25]</sup> TZP basis sets with small cores were used for geometry optimization and vibrational frequency analysis; all-electron basis sets were used for property calculations.<sup>[26]</sup> Scalar relativistic effects were taken into account by the ZORA to the Dirac equation.<sup>[27]</sup> The time-dependent DFT approach<sup>[28]</sup> was used to calculate the energies and oscillator strengths of excited states at the optimized ground state geometries. EPR parameters were calculated with inclusion of spin-orbit relativistic effects through the ZORA formalism.<sup>[29]</sup>

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